

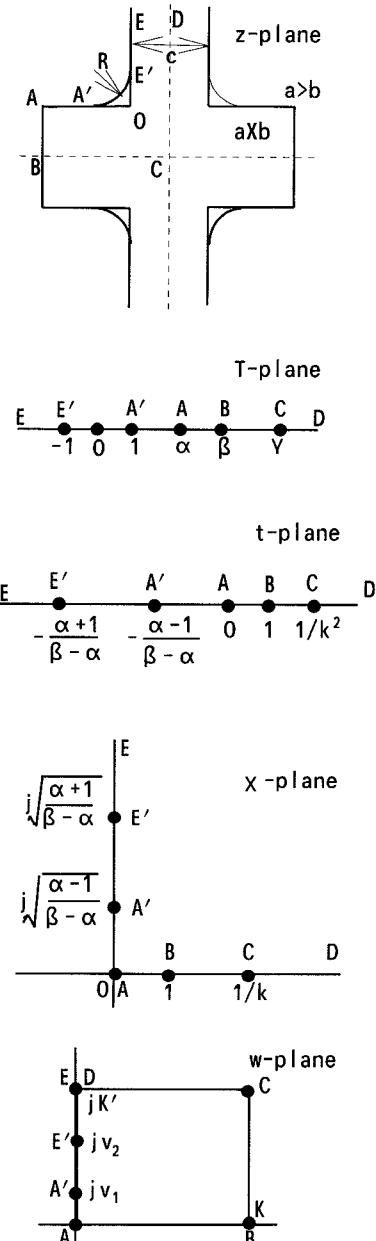
Jacobian Elliptic Function Solution of Characteristics of Rectangular Groove Wave-guide with Rounded Internal Corners

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ABSTRACT

To enhance the power carrying capacity and the efficiency of the rectangular groove wave guide, presented is such a guide with rounded internal corners, by means of Jacobian Elliptic Functions to transform the unbounded waveguide region onto the interior of a rectangle in w-plane with sides $2K(k) \times 2K'(k)$.



Rectangular or groove guides have sharp corners in the guide proper so are not low loss and high power carrying wave guides, theoretically. Proposed now is a groove rectangular waveguide with rounded internal angles, as shown in the figure, where R is radius of the circular rounded corner inside the guide. We map one quarter of this rectangular groove guide with rounded internal angles in the z-plane onto the upper half plane of the T-plane and proceed as Smythe dose (W.R.Smythe, Statics and Dynamic Electricity, § 4.28), by the following equation

$$\begin{aligned}
dz &= A_1 \frac{(T+1)^{1/2} + \lambda(T-1)^{1/2}}{\{(T-\alpha)(T-\beta)(T-\gamma)\}^{1/2}} dT \\
&= KA_1 \frac{(t+m)^{1/2} + \lambda(t+n)^{1/2}}{\{t(1-t)(1-k^2t)\}^{1/2}} dT
\end{aligned} \tag{1}$$

by a change of variable

$$t = \frac{T-\alpha}{\beta-\alpha}, \quad k^2 = \frac{\beta-\alpha}{\gamma-\alpha}, \quad m = \frac{\alpha+1}{\beta-\alpha}.$$

We have to find five constant $A_1, \lambda, \alpha, \beta$ and γ , from

the points $\infty, jR, -R, -\frac{a-c}{2}, -\frac{a-c}{2} - j\frac{b}{2}$ and $\frac{c}{2} - j\frac{b}{2}$. First, we let $v \rightarrow \infty$, then (1) gives

$$\frac{A_1(1+\lambda)T^{1/2}}{T^{3/2}} Tjd\theta = dU, \text{ integrating, we find}$$

$$A_1 = j \frac{c}{2\pi(1+\lambda)} \tag{2}$$

where c is the width of the groove. To make use of the other points, we make

$$t = \chi^2, \quad \chi = sn(w, k), \quad w = u + jv \tag{3}$$

and in the x and w planes points on the boundary of the z -plane are as shown, (1) becomes

$$dz = 2A_1 k \left\{ \sqrt{m + sn^2(w, k)} + \lambda \sqrt{n + sn^2(w, k)} \right\} dw \tag{4}$$

and the values of u_1 and u_2 can be found as follows

$$j \sqrt{\frac{\alpha+1}{\beta-\alpha}} sn(u + jv, k) = \frac{sd_1 + jcds_1 c_1}{1 - s_1^2 d^2} \tag{5}$$

where $s = sn(v, k)$, $s_1 = sn(u, k')$..., so $s(u, k) = 0$, $u=0$, and (5) gives

$$\sqrt{\frac{\alpha+1}{\beta-\alpha}} = \frac{c_1 s_1}{1 - s_1^2} = \frac{s_1(v_2)}{c_1(v_2)} = tn(v_2), \quad v_2 = tn^{-1} \left(\sqrt{\frac{\alpha+1}{\beta-\alpha}}, k' \right),$$

and similarly $v_1 = sn^{-1} \left(\sqrt{\frac{\alpha-1}{\beta-1}}, k' \right)$. We can now

integrate (5), starting from E' to C as follows

$$\begin{aligned}
-R(1+j) &= \int_{v_2}^{v_1} \frac{dz}{dw} dw \\
&= \int_{v_2}^{v_1} \left\{ \sqrt{m - \frac{sn^2(v_1, k)}{cn^2(v, k)}} + \sqrt{n - tn^2(v, k)} \right\} j 2A_1 k dv
\end{aligned}$$

$$-\frac{a-c}{2} + R = \int_{v_1}^{v_2} \left\{ \sqrt{m - tn^2 v} + \sqrt{n - tn^2 v} \right\} j A_1 k dv$$

$$-j \frac{b}{2} = \int_0^K \left\{ \sqrt{m + sn^2(u, k)} + \sqrt{n + sn^2(u, k)} \right\} 2A_1 k dv$$

$$\begin{aligned}
-\frac{a}{2} &= \int_K^{K+jK} \frac{dz}{dw} dw \\
&= \int_0^K \left\{ \sqrt{m + dn^2(y, k')} + \lambda \sqrt{n + dn^2(y, k')} \right\} j 2A_1 k dy
\end{aligned}$$

(by $w = K + jy$)

Solving these four equations numerically, we obtain α , β , γ and λ . We study the weighted wave equation of the rectangular groove guide with rounded internal angles

$$\frac{\partial^2 H_z}{\partial u^2} + \frac{\partial^2 H_z}{\partial v^2} + h^2 (\gamma^2 + k^2) H_z = 0, \quad h^2 = \left| \frac{dz}{dw} \right|^2$$

No exact solution of this equation can be founded, and we have to resort to numerical or approximate solutions. For example, in the zero order approximation we can neglect $sn^2(w, k)$ in (4) so h is a constant and this wave equation gives sinusoidal waves in a rectangular guide of dimensions of $2K(k) \times 2K'(k)$ in the w -plane as shown in the figure.

By transforming the boundary in the $\chi = \sqrt{t}$ plane by the Jacobian elliptic function, we transform the unbounded guide region in the z -plane into the bounded region of the interior of a rectangle in the w -plane, we are then in a better position to solve the weighted wave equation by numerical or approximate method.